

## SUB-SOLAR $\Delta^{17}\text{O}$ VALUES IN PARENT CLOUD CORE INFALL MODELS.

J. R. Lyons<sup>1</sup> and J.-E. Lee<sup>2</sup>,  
<sup>1</sup>School of Earth & Space Exploration, ASU, jimlyons@asu.edu, <sup>2</sup>Astronomy & Space Science, School of Space Research, Kyung Hee University, Korea.

**Introduction:** Recent oxygen isotope measurements of grossite-rich CAIs with low initial  $^{26}\text{Al}/^{27}\text{Al}$  ratio from CH chondrites exhibit  $\Delta^{17}\text{O}$  values as low as -40 ‰ [1]. Assuming that CAIs with low initial  $^{26}\text{Al}/^{27}\text{Al}$  predate CAIs with the canonical  $^{26}\text{Al}/^{27}\text{Al}$  ratio ( $\sim 5 \times 10^{-5}$ ), infall-disk models suggest that these very early CAIs may have formed ~20 Kyr – 30 Kyr after the initiation of collapse of the parent cloud core that formed the solar system [2]. This very short timescale is most consistent with inheritance of low  $\Delta^{17}\text{O}$  values derived from CO self-shielding in the parent cloud [3], [4]. Here, we explore the dust/gas fractionation during the infall stage, and the implications for infall material entering the growing disk. Recently, measurements of  $\Delta^{17}\text{O}$  layering in CAI spherules have been interpreted as evidence for inheritance of self-shielding effects from the parent cloud [5].

**Transport timescales in accretion disk:** A likely explanation for the very low  $\Delta^{17}\text{O}$  values reported in [1] is rapid transport of self-shielded CO along the UV-active disk surface. Radial transport timescales in an accretion disk, such as that used for modeling CO

$$\text{self-shielding} [6], \text{ are calculated as } t_r = \frac{R^2}{ac_s H} \text{ for}$$

heliocentric distance  $R$ , sound speed  $c_s$ , and vertical scale height  $H$ . For CO self-shielding in the disk surface, only  $\alpha \sim 0.1$  yields timescales short enough to explain the low  $^{26}\text{Al}/^{27}\text{Al}$  CAI data. The UV-active surface can, indeed, have  $\alpha$  as high as 0.1, so this is a seemingly plausible scenario. However, during the early stages of infall, the disk does not have an accretion disk structure, as assumed in [6], and implicit in Figure 1. We therefore consider delivery of self-shielded CO in the cloud core infall models.

**Cloud core infall model:** We use the cloud core model of Lee et al. [4], and explore the role of gas/dust fractionation in this model. Infall is modeled by densification of a Bonnor-Ebert sphere, followed by the start of gravitational collapse, which defines  $t = 0$  in the model. We define  $\beta$  as the fraction of  $\text{H}_2\text{O}$  remaining in the infall gas at the inner boundary of 125 AU of the 1-D spherical model. Dust/gas fractionation is a material fractionation process that segregates CO gas with  $\Delta^{17}\text{O} \ll \text{solar}$  in the top of the growing disk from  $\text{H}_2\text{O}$  ice with  $\Delta^{17}\text{O} \gg \text{solar}$ . CO self-shielding and formation of  $\text{H}_2\text{O}$  from  $^{17}\text{O}$  and  $^{18}\text{O}$ -rich O atoms create the large isotopic difference in these two O res-

ervoirs. The process of dust/gas fractionation requires a physical separation of the two reservoirs by, e.g., gravitational settling of ice-coated dust grains. The densification phase of the infall model affects the ability of FUV photons to penetrate the cloud core, and results in CO self-shielding signatures that vary greatly with the magnitude of the FUV radiation field incident on the cloud core [4].

Figure 2 a-f show the resulting  $\Delta^{17}\text{O}$  values for dust/gas fractionation infall material at the inner boundary of the infall model.  $G_0$  is a measure of the radiation field and is a multiplicative factor amplifying the local interstellar medium (ISM) FUV field. The required dust/gas fractionation factor,  $\beta$ , varies greatly with  $G_0$ , with no solution for  $G_0 = 1$  and  $G_0 = 10^3$ , to  $\beta = 0.9$  for  $G_0 = 100$  and  $G_0 = 10^5$ .

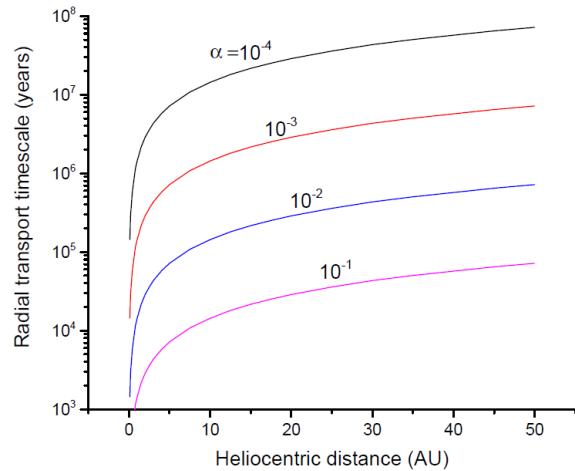


Figure 1. Radial transport timescales in an accretion disk, assuming a turbulent viscosity parameter,  $\alpha$ . (Figure from [7]).

**Conclusions:** We find that for some radiation fields only a slight dust/gas fractionation is sufficient to explain the CH CAI  $\Delta^{17}\text{O}$  data in [1]. A full evaluation of this requires coupling of our infall model into a growing disk model [2]. This is our next step.

**References:** [1] Krot A. N. et al. (2019) LPSC 50th, abstract #1230. [2] Pignatale F. C. et al. (2018) *ApJ Lett.*, 867, L23. [3] Yurimoto H. and Kuramoto K. (2004) *Science*, 305, 1763. [4] Lee J.-E. et al. (2008) *Meteor. & Planet. Sci.*, 43, 1351. [5] Simon J. et al. (2019) *Ap. J.*, 884, L29. [6] Lyons J. and Young E. (2005) *Nature*, 435, 317. [7] Krot A. N. et al. (submitted).

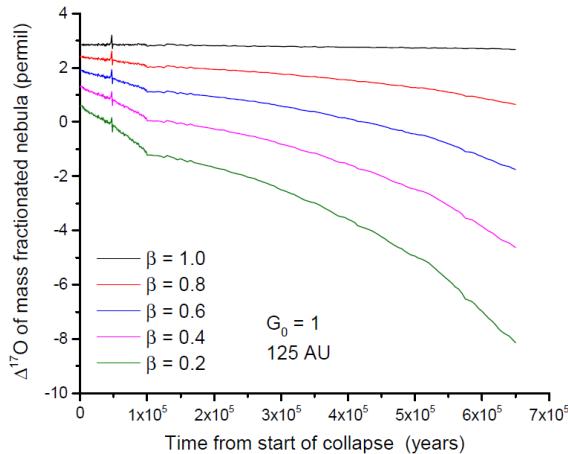


Figure 2a.  $\Delta^{17}\text{O}$  for dust/gas fractionated material at the inner boundary of a cloud core infall model.  $\beta$  is the fraction of  $\text{H}_2\text{O}$  remaining with the gas at the inner boundary.  $G_0 = 1$ , which corresponds to the ISM FUV field. The CH CAIs require  $\Delta^{17}\text{O} \sim 12\text{\textperthousand}$  or less at a few  $\times 10^4$  years.

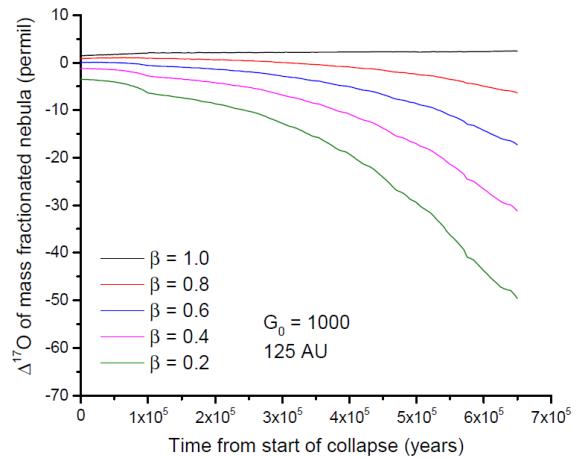
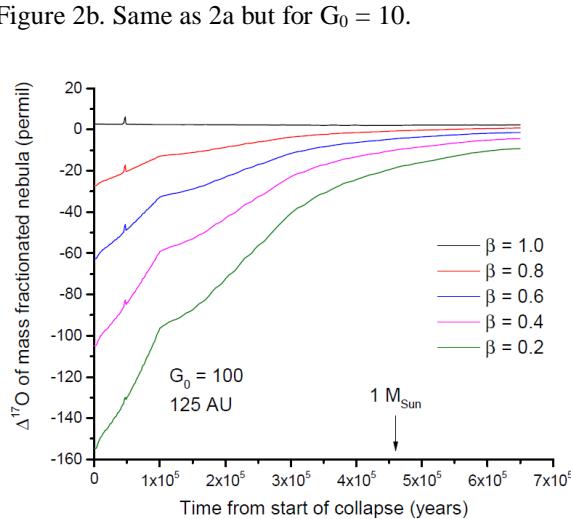
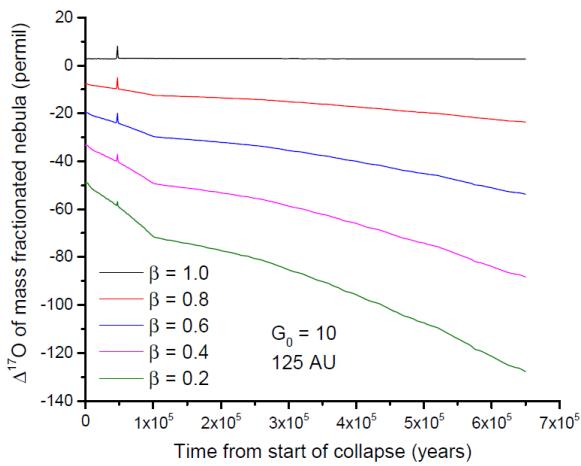


Figure 2d. Same as 2a but for  $G_0 = 1000$ . No value of  $\beta$  is consistent with the CH CAI data in [1].

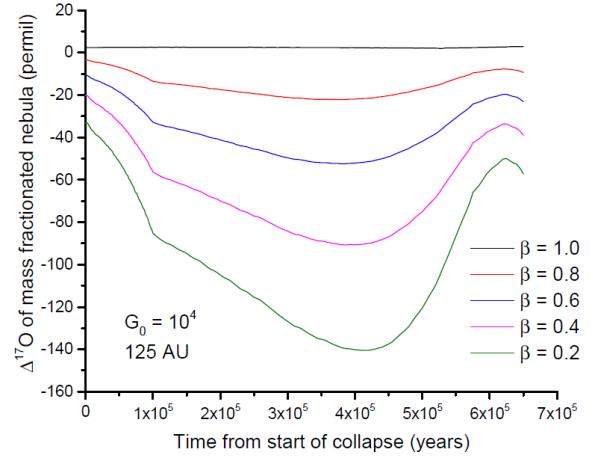


Figure 2e. Same as 2a but for  $G_0 = 1 \times 10^4$ .

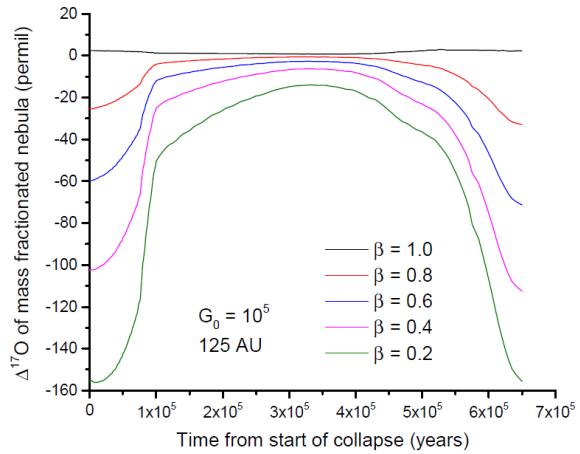


Figure 2f. Same as 2a but for  $G_0 = 1 \times 10^5$ .